

The proton mass from first principles: lattice QCD at the physical point

Martha Constantinou



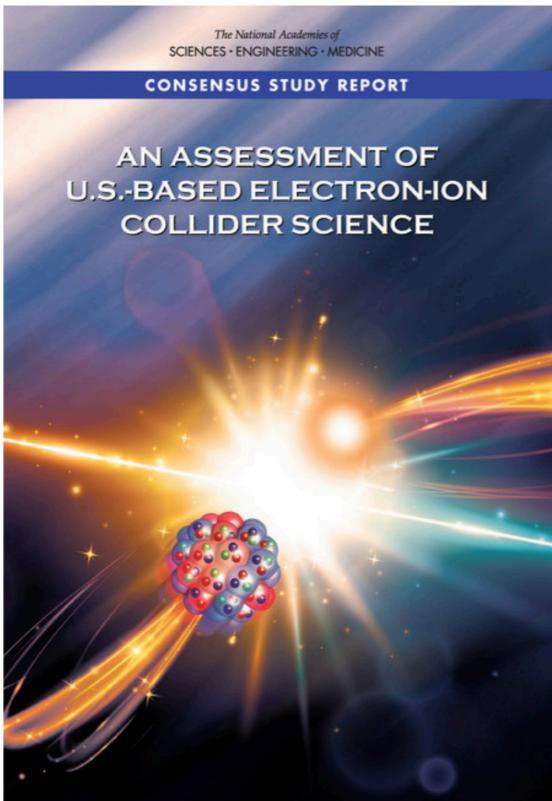
In collaboration with: **C. Alexandrou, K. Hadjiyiannakou**



3rd Proton Mass Workshop; Origin and Perspective

January 14, 2021

Proton Mass



Main Pillar of NAS Assessment report for EIC

Finding 1: An EIC can uniquely address three profound questions about nucleons—neutrons and protons—and how they are assembled to form the nuclei of atoms:

- How does the mass of the nucleon arise?
- How does the spin of the nucleon arise?
- What are the emergent properties of dense systems of gluons?

Lattice QCD can provide valuable input in understanding the proton mass decomposition from *first principles*

... before experimental EIC data are available

Proton Mass Decomposition

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- ★ Starting point is the symmetric EMT, relevant in the fwd limit
(Non-symm. part vanishes by e.o.m. of quarks and gluons)

$$T_{sym}^{\mu\nu} = \frac{1}{4}\bar{\psi}i\overleftrightarrow{D}^{(\mu}\gamma^{\nu)}\psi - F^{\mu\alpha}F_\alpha^\nu + \frac{g^{\mu\nu}}{4}F^{\alpha\beta}F_{\alpha\beta}$$

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$$\langle T^{\mu\nu} \rangle = 2P^\mu P^\nu$$

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- ★ In rest-frame, the mass is related to the matrix elements of EMT

$$\frac{\langle T_\mu^\mu \rangle}{\langle N|N \rangle} = M, \quad \frac{\langle T^{00} \rangle}{\langle N|N \rangle} = M$$

Proton Mass Decomposition

Based on sum rules
(not unique)

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see, e.g., [M. Shifman et al., Phys. Lett. 78B (1978); D. Kharzeev, Proc. Int. Sch. Phys. Fermi 130 (1996)]

★ Decomposition of T^{00} in trace and traceless parts in rest frame

[X.D. Ji, Phys. Rev. Lett. 74, 1071 (1995); X. D. Ji, Phys. Rev. D 52, 271 (1995)]

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[C. Lorce', Eur. Phys. J. C78 (2018) 2, arXiv:1706.05853]

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Once the EMT is decomposed into components,
renormalization of the latter is necessary

Ji's Decomposition

$$\frac{\langle T^{00} \rangle}{\langle N | N \rangle} = M$$

Ji' Decomposition

[X.D. Ji, Phys. Rev. Lett. 74, 1071 (1995); X. D. Ji, Phys. Rev. D 52, 271 (1995)]

- ★ Traceless ($\bar{T}^{\mu\nu}$) & trace ($\hat{T}^{\mu\nu}$) parts of EMT:

$$T^{\mu\nu} = T_q^{\mu\nu} + T_g^{\mu\nu}, \quad T_{q,g}^{\mu\nu} = \bar{T}_{q,g}^{\mu\nu} + \hat{T}_{q,g}^{\mu\nu}$$

- ★ Trace of EMT: $\hat{T}^{\mu\nu} = \frac{1}{4}g^{\mu\nu} \left[(1 + \gamma_m) \bar{\psi} m \psi + \frac{\beta(g)}{2g} F^2 \right]$

- ★ $\langle T^{00} \rangle$ has for contributions from $\langle \hat{T}_q^{00} \rangle, \langle \bar{T}_q^{00} \rangle, \langle \hat{T}_g^{00} \rangle, \langle \bar{T}_g^{00} \rangle$

- ★ energy density component gives a decomposition for the mass:

$$M = \frac{\langle N | T^{00} | N \rangle}{\langle N | N \rangle} = M_m + M_q + M_g + M_a$$

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quark mass $\langle \hat{T}_q^{00} \rangle$
quark energy $\langle \bar{T}_q^{00} \rangle$
gluon energy $\langle \bar{T}_g^{00} \rangle$
trace anomaly $\langle \hat{T}_g^{00} \rangle$

Question:

“Can lattice calculate the mass distribution in the nucleon?”

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Answer:

Components associated with operators calculable in lattice QCD

★ Quark mass

$$M_m = \sum_q \sigma_q$$

★ Quark energy

$$M_q = \frac{3}{4} \left(M \sum_q \langle x \rangle_q - \sum_q \sigma_q \right)$$

★ Gluon energy

$$M_g = \frac{3}{4} M \langle x \rangle_g$$

★ Trace anomaly

$$M_a = \frac{\gamma_m}{4} \sum_q \sigma_q - \frac{\beta(g)}{4g} (E^2 + B^2)$$

σ_q : sigma-terms

$\langle x \rangle_q$: Quark momentum fraction

$\langle x \rangle_g$: Gluon momentum fraction

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Results at the physical point

[C. Alexandrou et al., PRD 102, 054517 (2020), arXiv:1909.00485]

See Alexandrou's talk

[C. Alexandrou et al., PRD 101, 094513 (2020), arXiv:2003.08486]

σ_q : sigma-terms

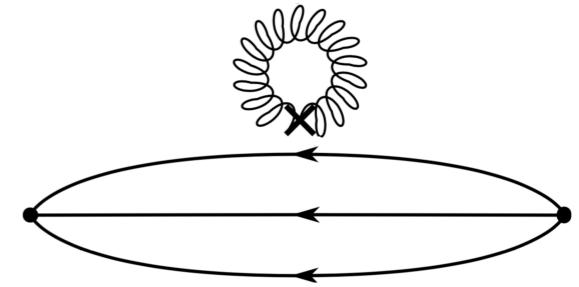
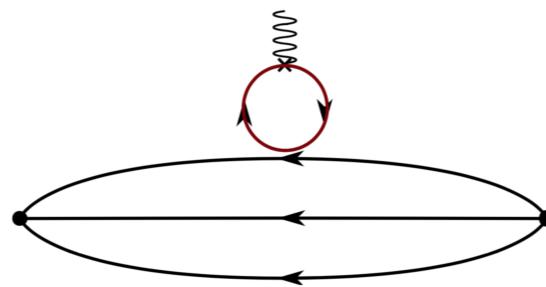
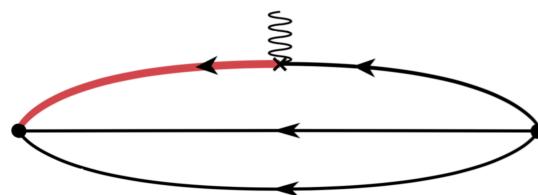
$\langle x \rangle_q$: Quark momentum fraction

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Results @ physical pion mass

$\overline{MS}(2\text{GeV})$

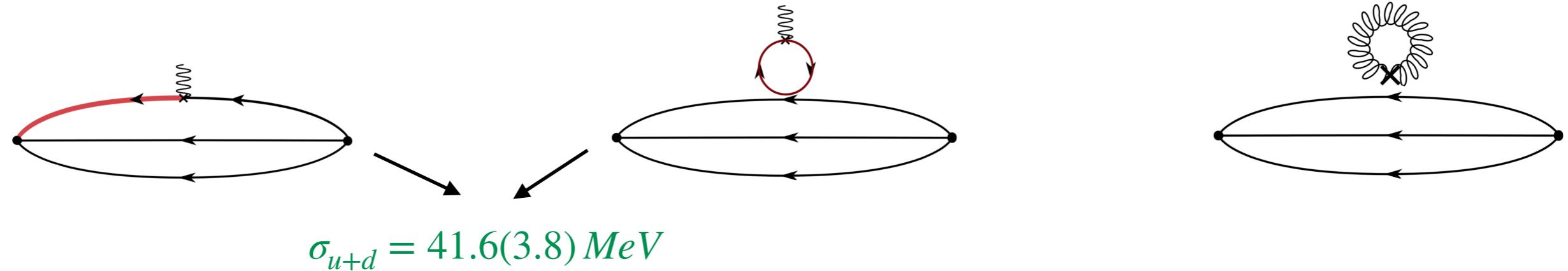
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Results @ physical pion mass

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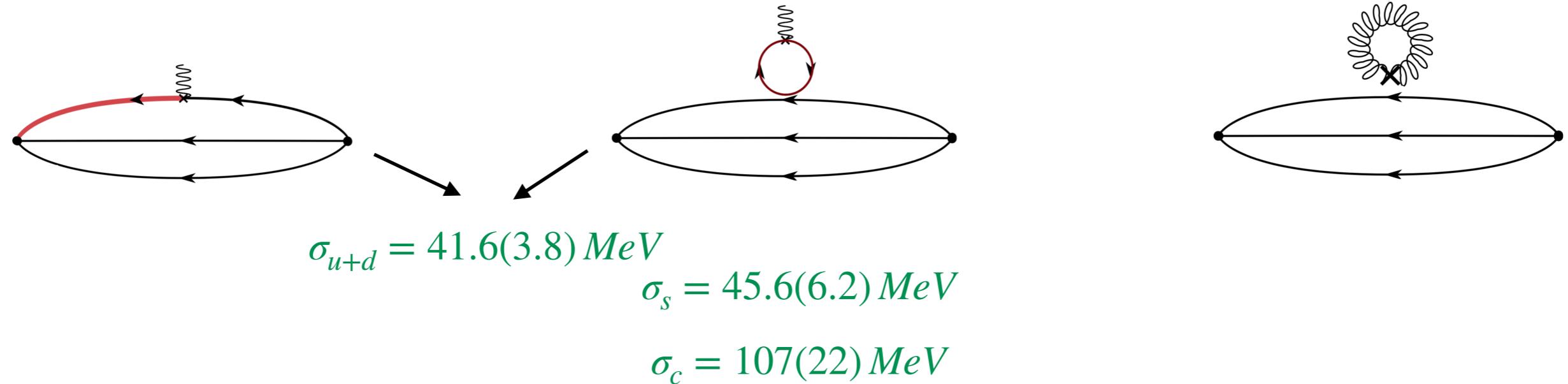
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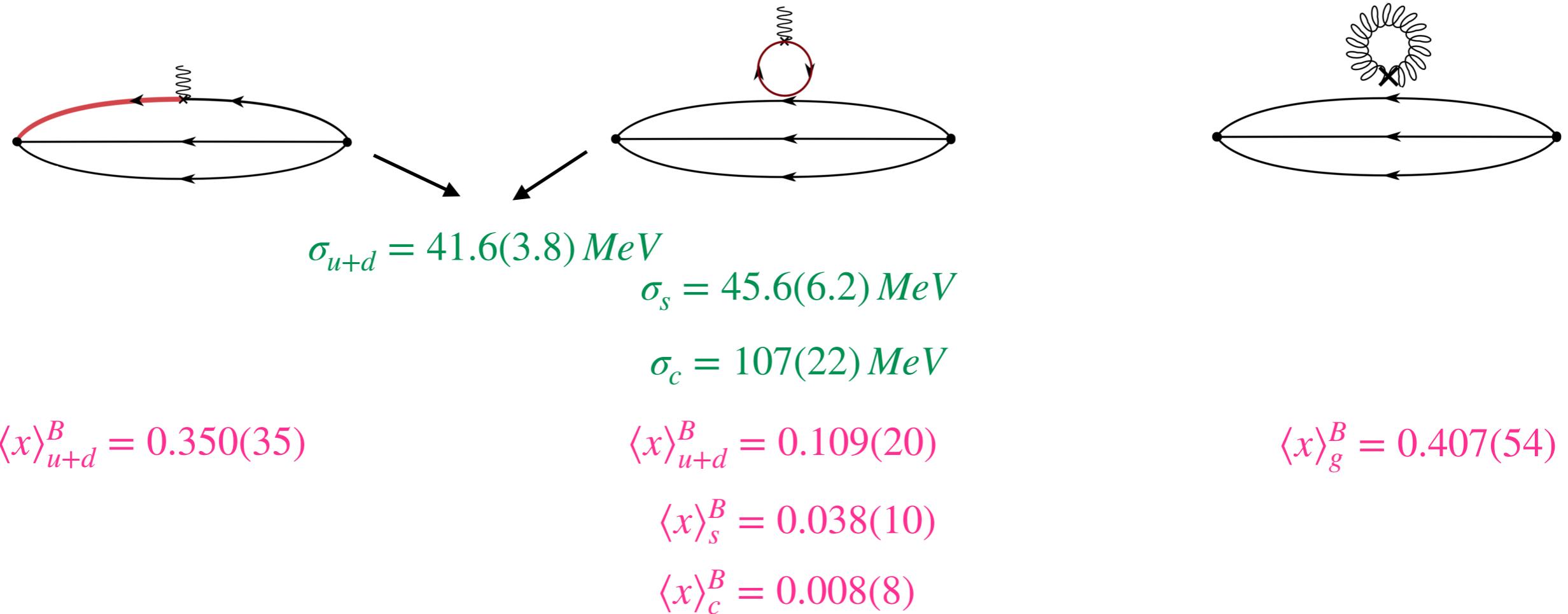
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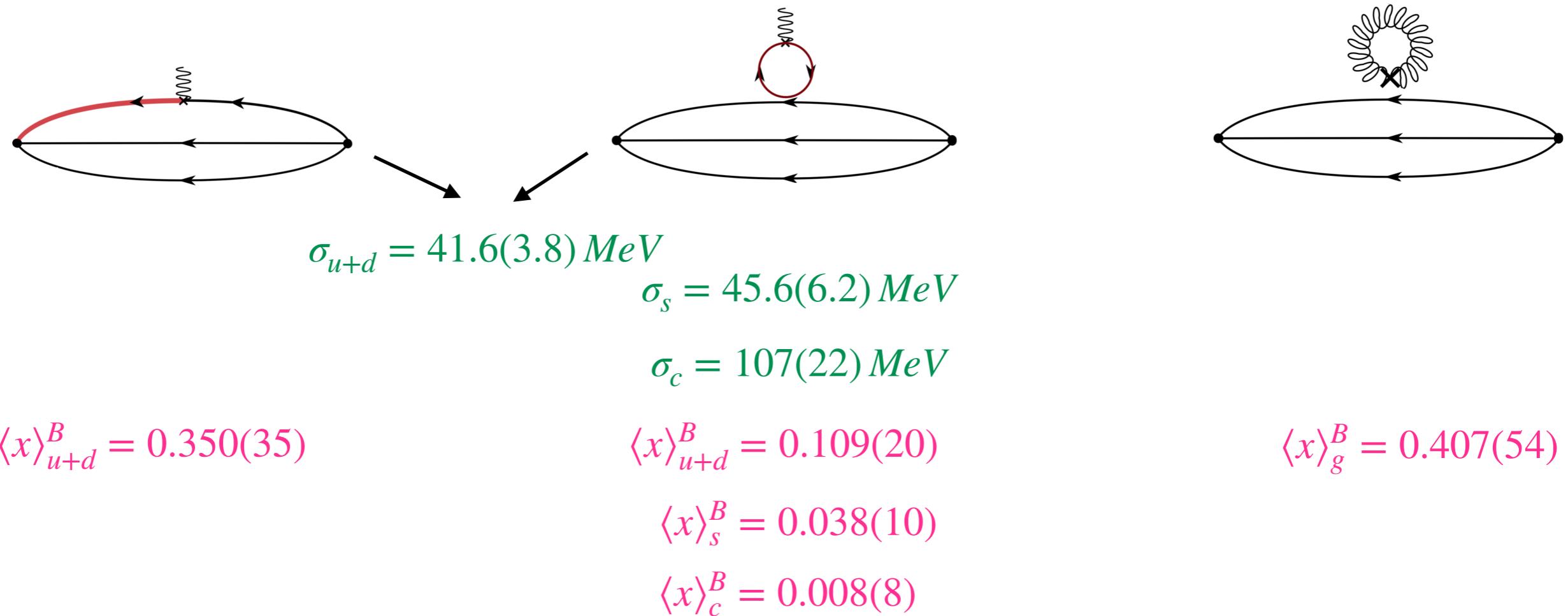
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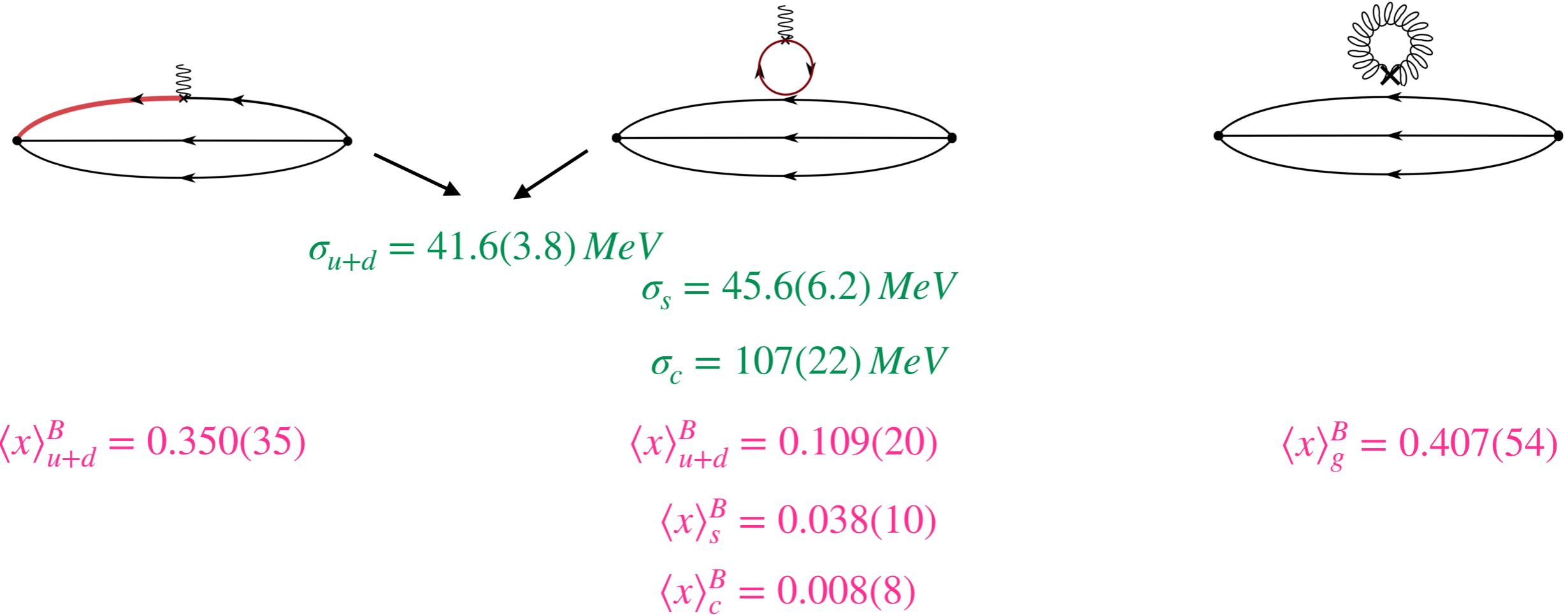
★ Mixing between quark and gluon contributions to $\langle x \rangle$

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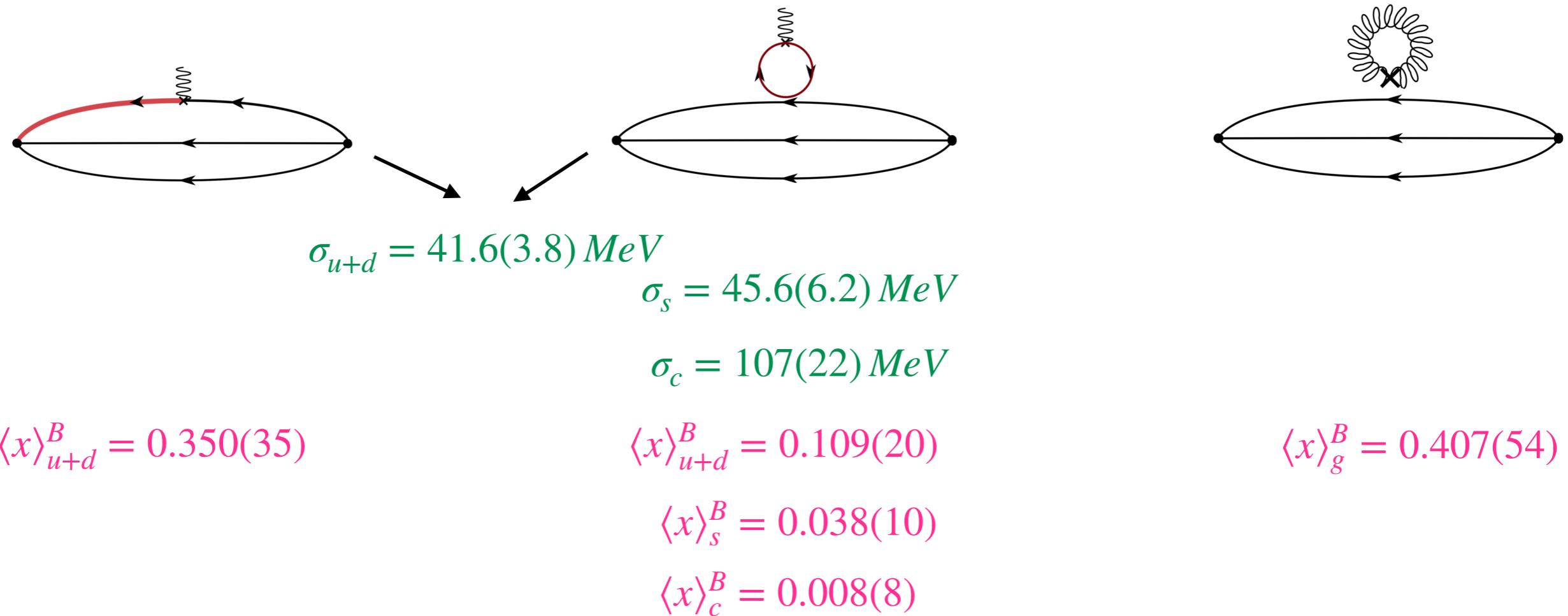
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$$\langle x \rangle_u = 0.359(30) \quad \langle x \rangle_d = 0.188(19) \quad \langle x \rangle_s = 0.052(12) \quad \langle x \rangle_c = 0.019(9) \quad \langle x \rangle_g = 0.427(92)$$

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Momentum sum rule satisfied!

Ji's Decomposition

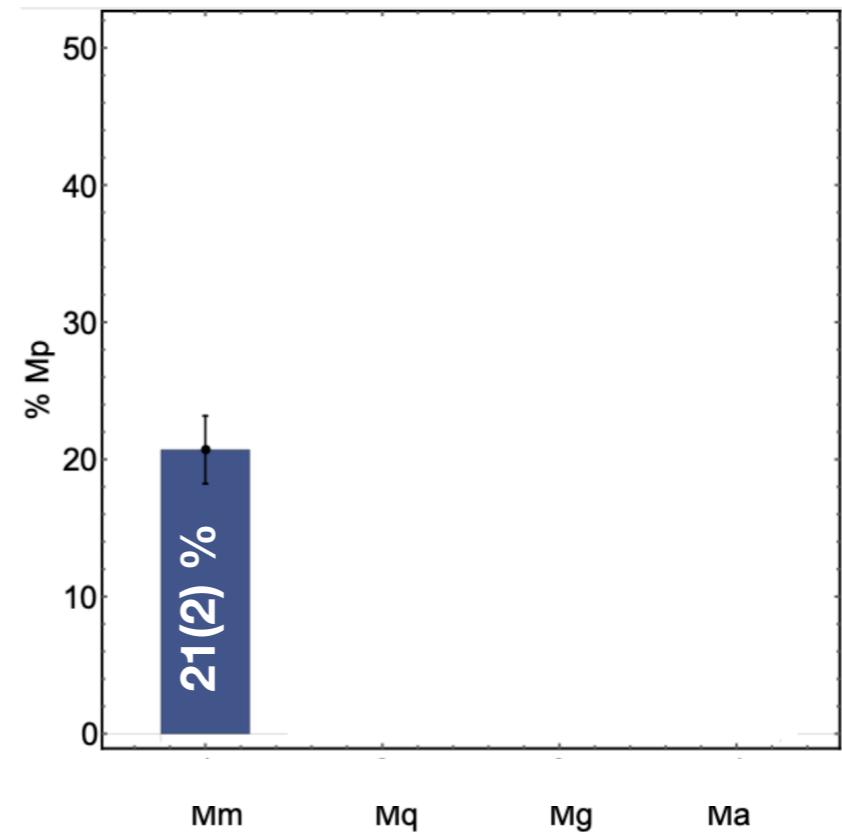
Proton Mass Budget

- ★ Available contributions:

Ji's Decomposition

Proton Mass Budget

- ★ Available contributions:
 - quark mass (σ -terms)

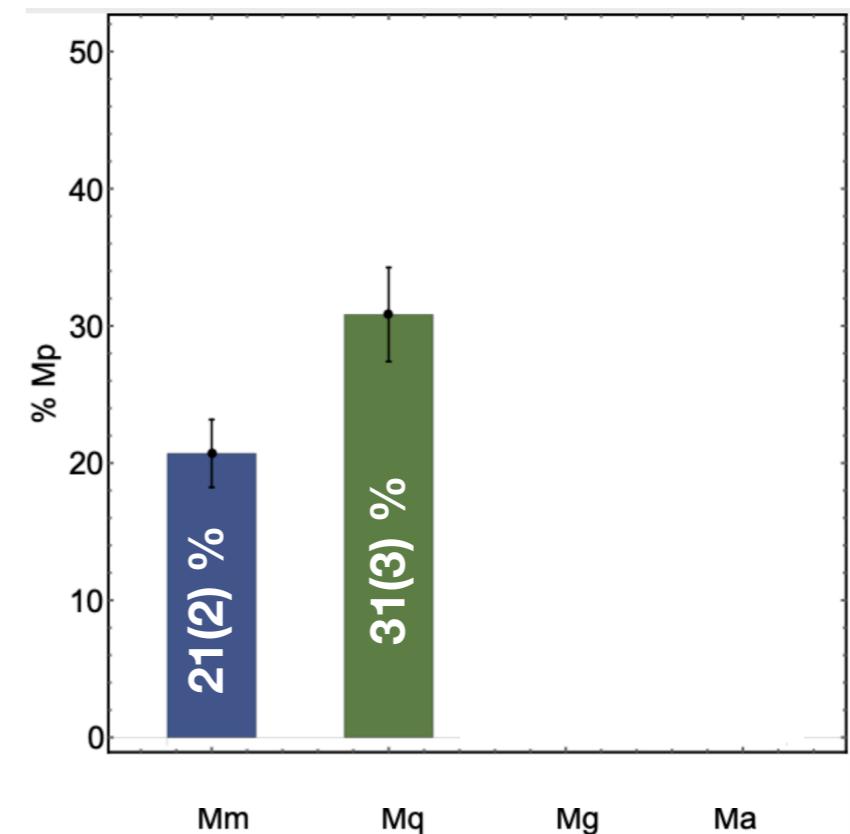


Ji's Decomposition

Proton Mass Budget

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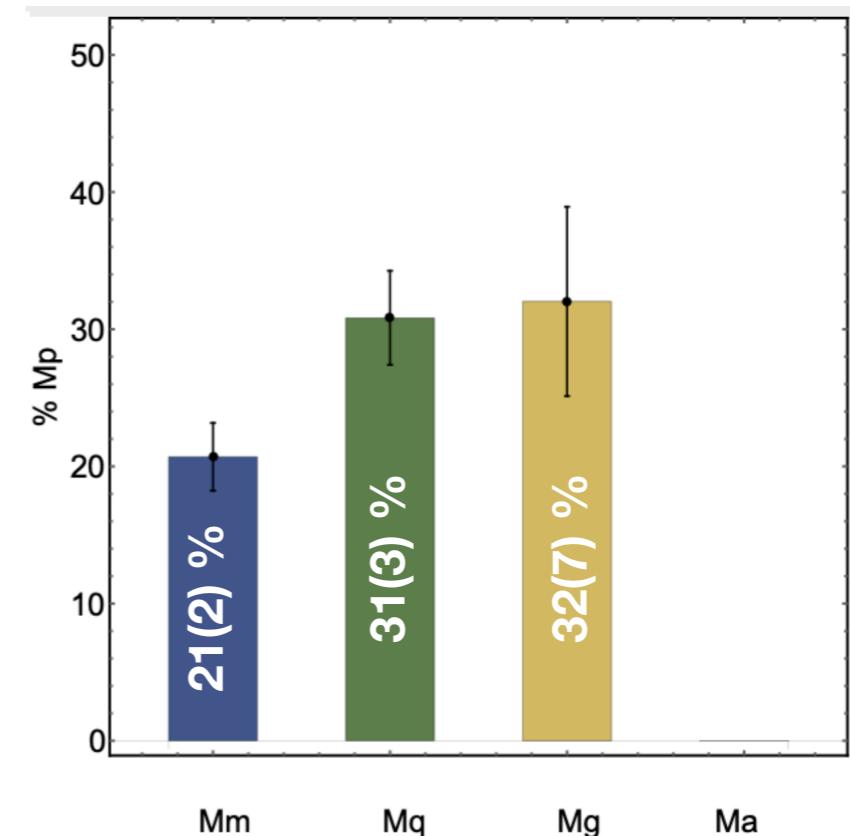


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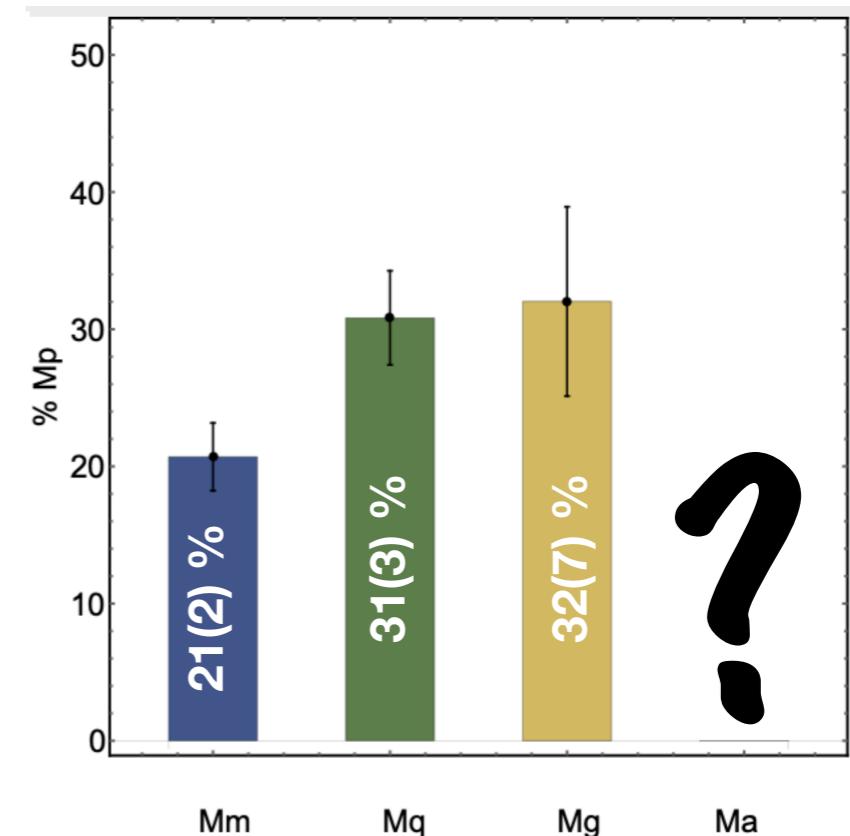


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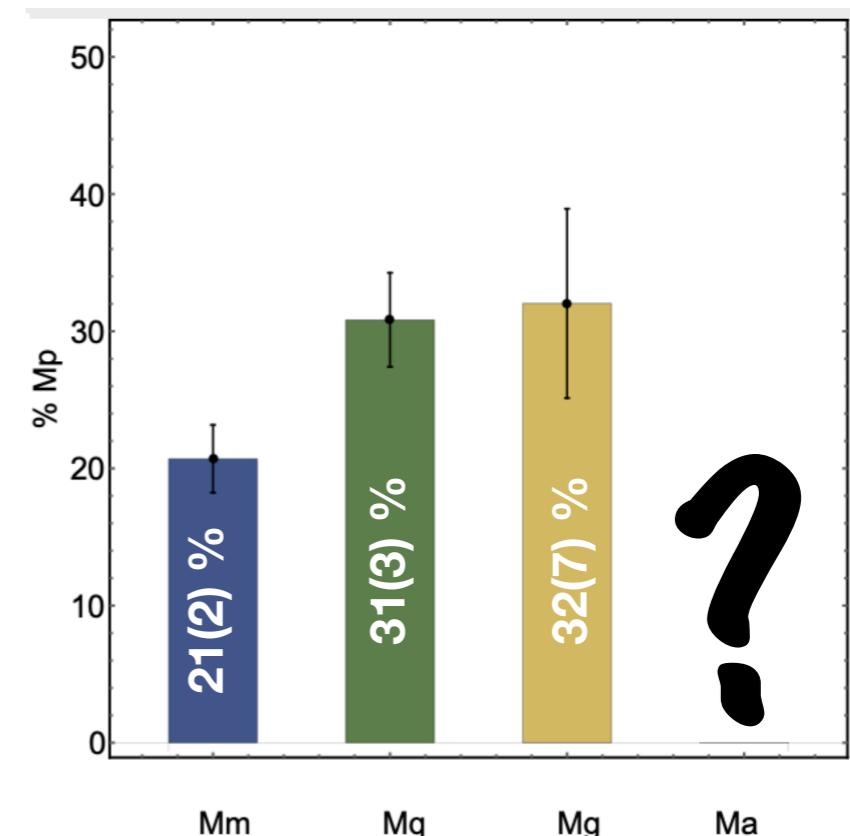
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Currently not available



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Theoretical & Technical challenges

- ★ Disconnected contribution (signal-to-noise ratio suppressed)
- ★ Presence of mixing with operators that are BRST variations and that vanish by the e.o.m.
(Full EMT: 10 renormalization functions)
(Trace: 3-operator mixing under renormalization in the continuum)

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Answer:

- ★ Direct calculation of trace anomaly not available
- ★ Possibility to access trace anomaly indirectly from sum rules

$$M_a = \frac{M}{4} - \sum_q \frac{\sigma_q}{4} \quad M_a = M - \sum_{i=m,q,g} M_i$$

Ji's Decomposition

Proton Mass Budget

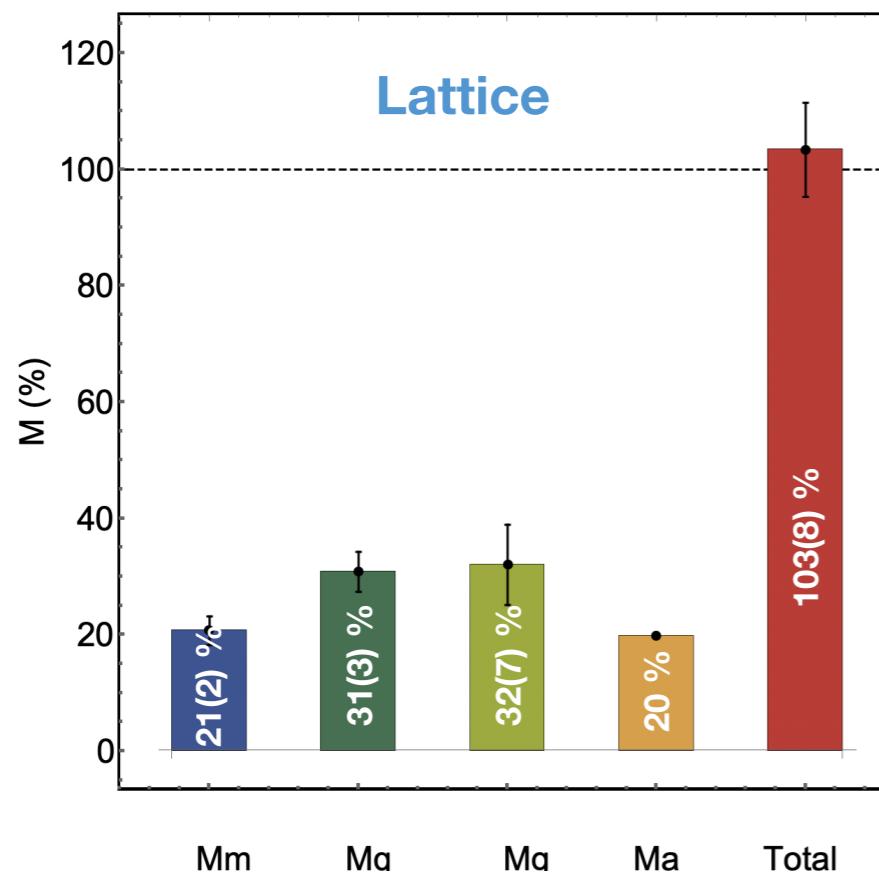
Ji's Decomposition

Proton Mass Budget

Approach A

$$M_a = \frac{M_p}{4} - \sum_q \frac{\sigma_q}{4} \sim 19.83(0.07) \%$$

$$M_p = M_m + M_q + M_g + M_a = 103.39(8.09) \%$$



Ji's Decomposition

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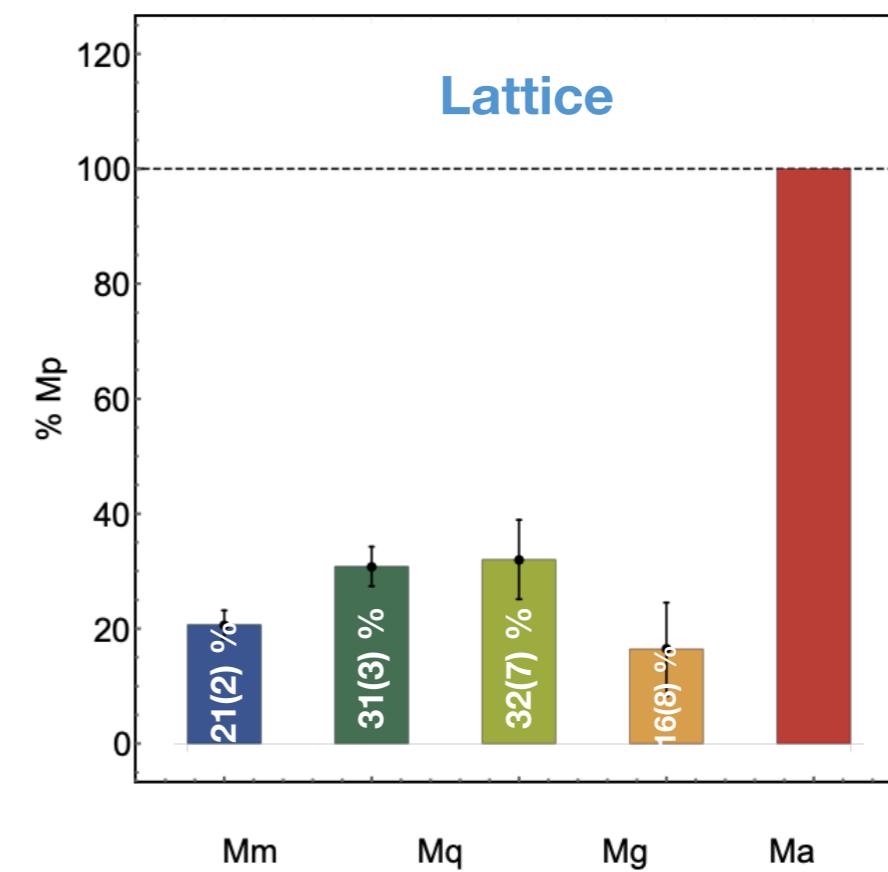
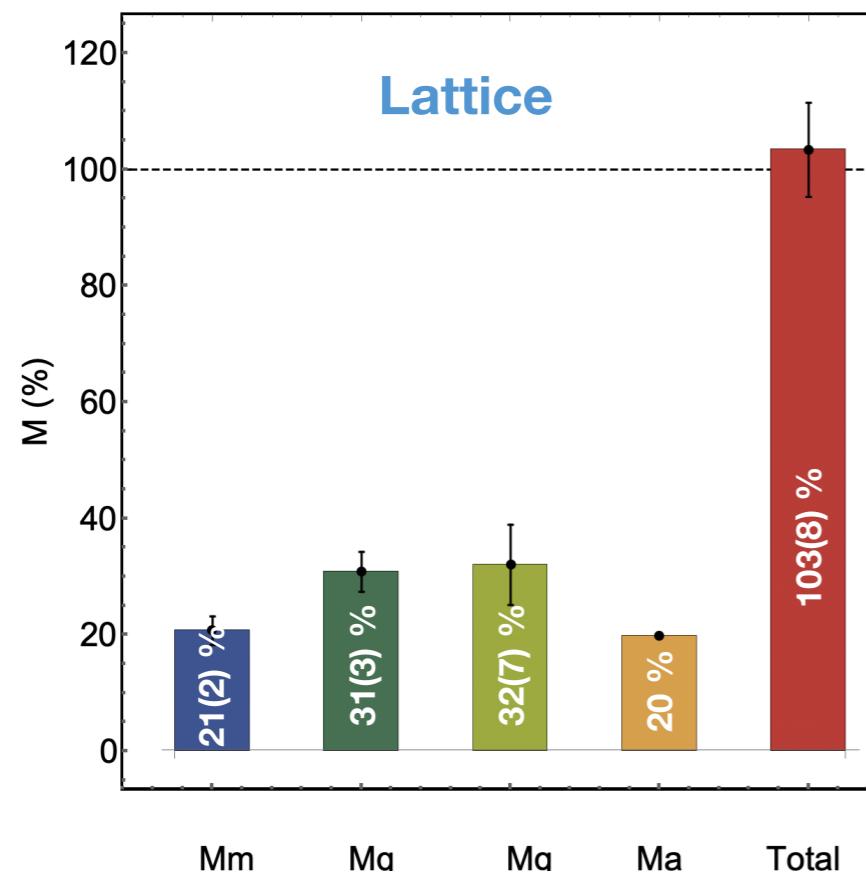
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Approach B

$$M_a = M_p - \sum_{i=m,q,g} M_i \sim 16.45(8.09) \%$$

M_p : sum rule enforced



Ji's Decomposition

Proton Mass Budget

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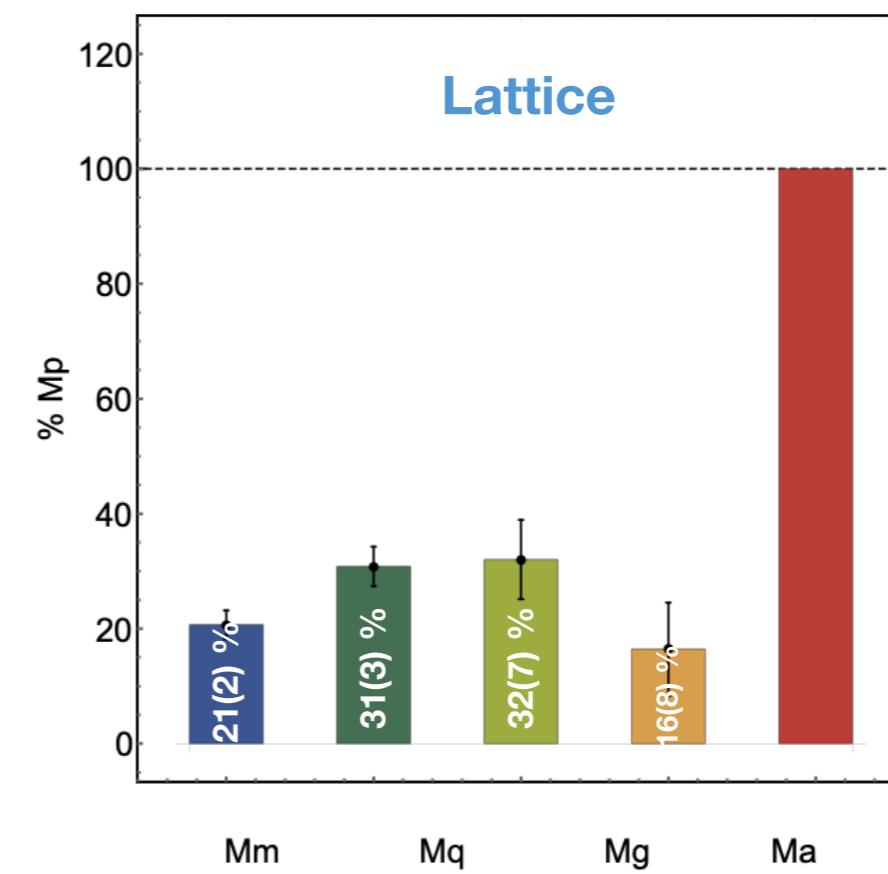
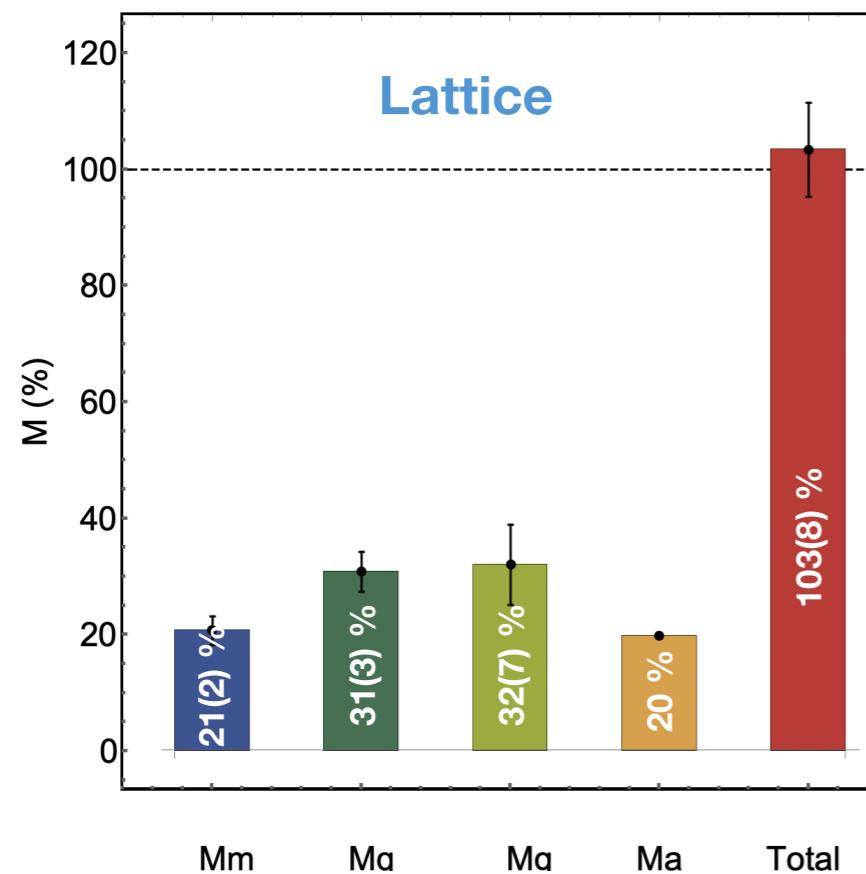
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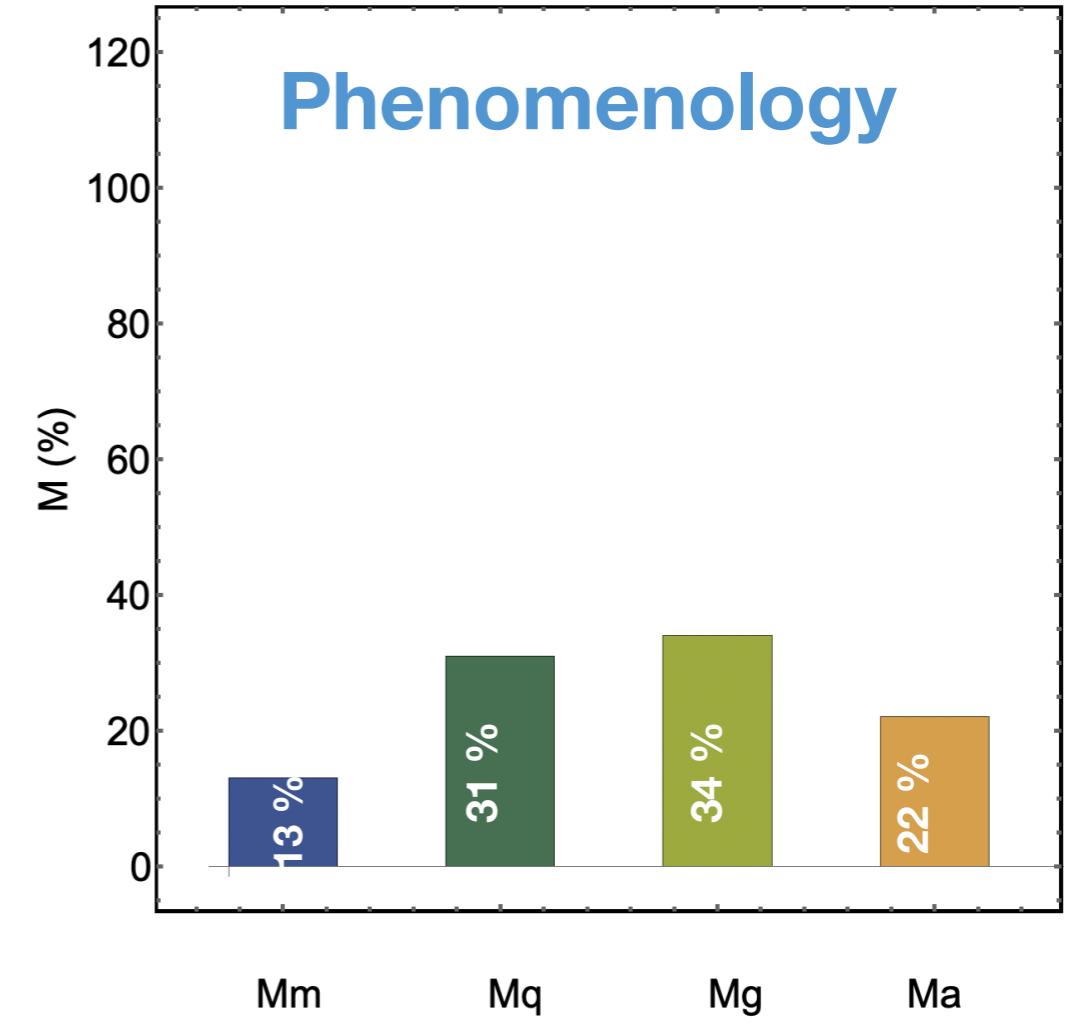
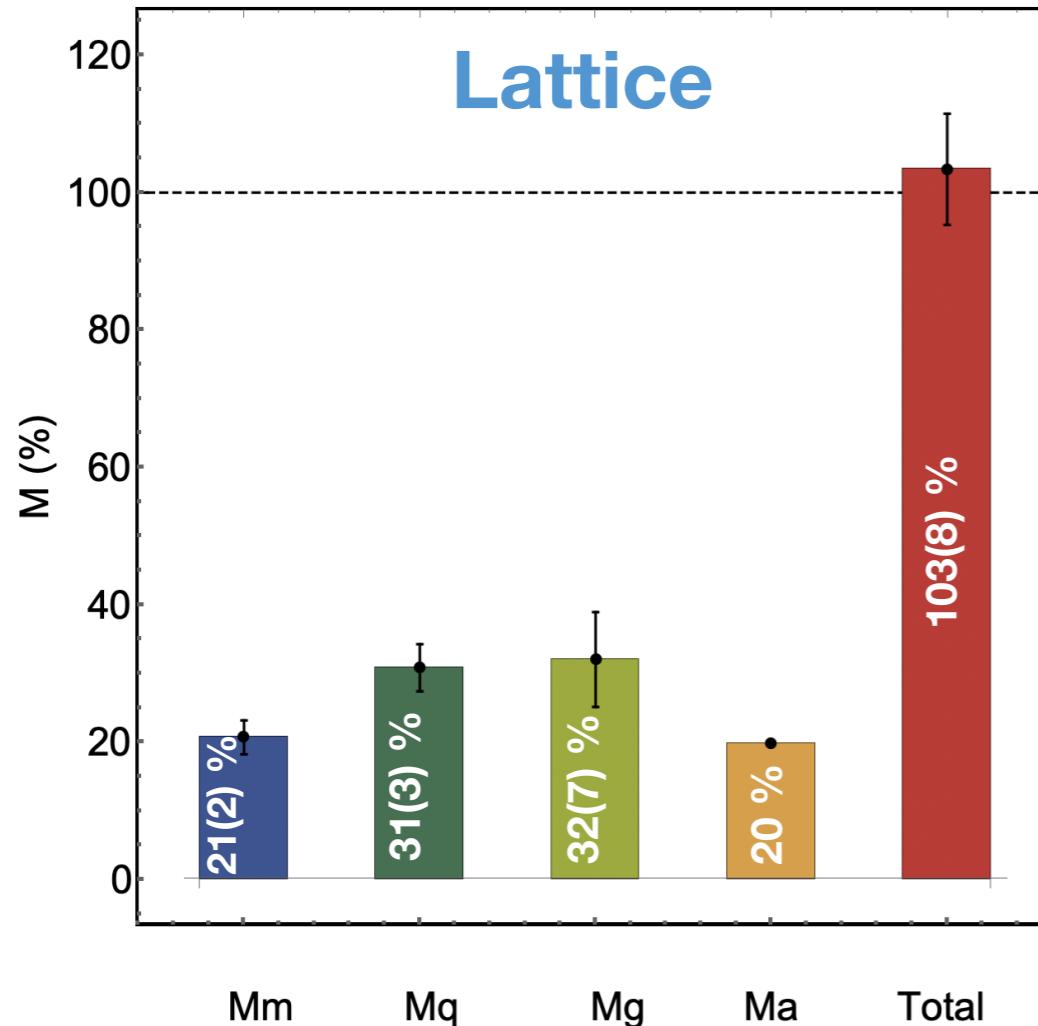
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M_p : sum rule enforced



- ★ **M_a compatible but different systematic uncertainties**
- ★ **Uncertainties of trace anomaly term depend on the sum rule**

Ji's Decomposition



- ★ Lattice and pheno data give similar picture
- ★ The tension in the sigma terms affects M_m
- ★ Contributions are of similar order

[C. Lorce', EPJ. C78 (2018) 2]
[L. Harland-Lang et al., EPJ. C 75 (2015)]
[M. Hoferichter et al., PRL 115 (2015)]

Lorcé's Decompositions

$$\frac{\langle T^{00} \rangle}{\langle N | N \rangle} = M$$

Decomposition with Pressure effects

[C. Lorce', Eur. Phys. J. C78 (2018) 2, arXiv:1706.05853]

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- ★ Matrix elements of each component
in the fwd limit:

$$\langle N | T_{q,g}^{\mu\nu}(0) | N \rangle = 2P^\mu P^\nu A_{q,g}(0) + 2M^2 g^{\mu\nu} \bar{C}_{q,g}(0)$$

$$(\langle T^{\mu\nu} \rangle = 2P^\mu P^\nu)$$

Sum rules

$$A_q(0) + A_g(0) = 1$$

$$\bar{C}_q(0) + \bar{C}_g(0) = 0$$

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- ★ The EMT FFs can be related to the components of Ji's picture:

$$A_q(0) = \sum_f \langle x \rangle_f, \quad A_g(0) = \langle x \rangle_g, \quad \bar{C}_q(0) = (1 + \gamma_m) \frac{\sigma_{u+d+s+c}}{4M} - \frac{\langle x \rangle_{u+d+s+c}}{4}, \quad \bar{C}_g(0) = -\bar{C}_q(0)$$

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$$T_g^{\mu\nu} = -F^{\mu\alpha} F_\alpha^\nu + \frac{g^{\mu\nu}}{4} F^{\alpha\beta} F_{\alpha\beta}$$

- ★ Matrix elements of each component in the fwd limit:

$$\langle N | T_{q,g}^{\mu\nu}(0) | N \rangle = 2P^\mu P^\nu A_{q,g}(0) + 2M^2 g^{\mu\nu} \bar{C}_{q,g}(0)$$

($\langle T^{\mu\nu} \rangle = 2P^\mu P^\nu$)
Sum rules

$$A_q(0) + A_g(0) = 1$$

$$\bar{C}_q(0) + \bar{C}_g(0) = 0$$

- ★ The EMT FFs can be related to the components of Ji's picture:

$$A_q(0) = \sum_f \langle x \rangle_f, \quad A_g(0) = \langle x \rangle_g, \quad \bar{C}_q(0) = (1 + \gamma_m) \frac{\sigma_{u+d+s+c}}{4M} - \frac{\langle x \rangle_{u+d+s+c}}{4}, \quad \bar{C}_g(0) = -\bar{C}_q(0)$$

$$A_q(0) + A_g(0) = 1 \quad \checkmark$$

Decomposition with Pressure effects

[C. Lorce', Eur. Phys. J. C78 (2018) 2, arXiv:1706.05853]

- ★ Thermodynamic potentials using the energy component $T_{q,g}^{00}$:
(effective coupled two-fluid picture, combinations of internal energies and pressure-volume works)

U : internal energy (kinetic & potential)

W : pressure-volume work)

Sum rules

$$M = U_q + U_g$$

$$W_q + W_g = 0$$

$$U_{q,g} = M \left[A_{q,g}(0) + \bar{C}_{q,g}(0) \right], \quad W_{q,g} = -M \bar{C}_{q,g}(0)$$

energy density and pressure kept separately

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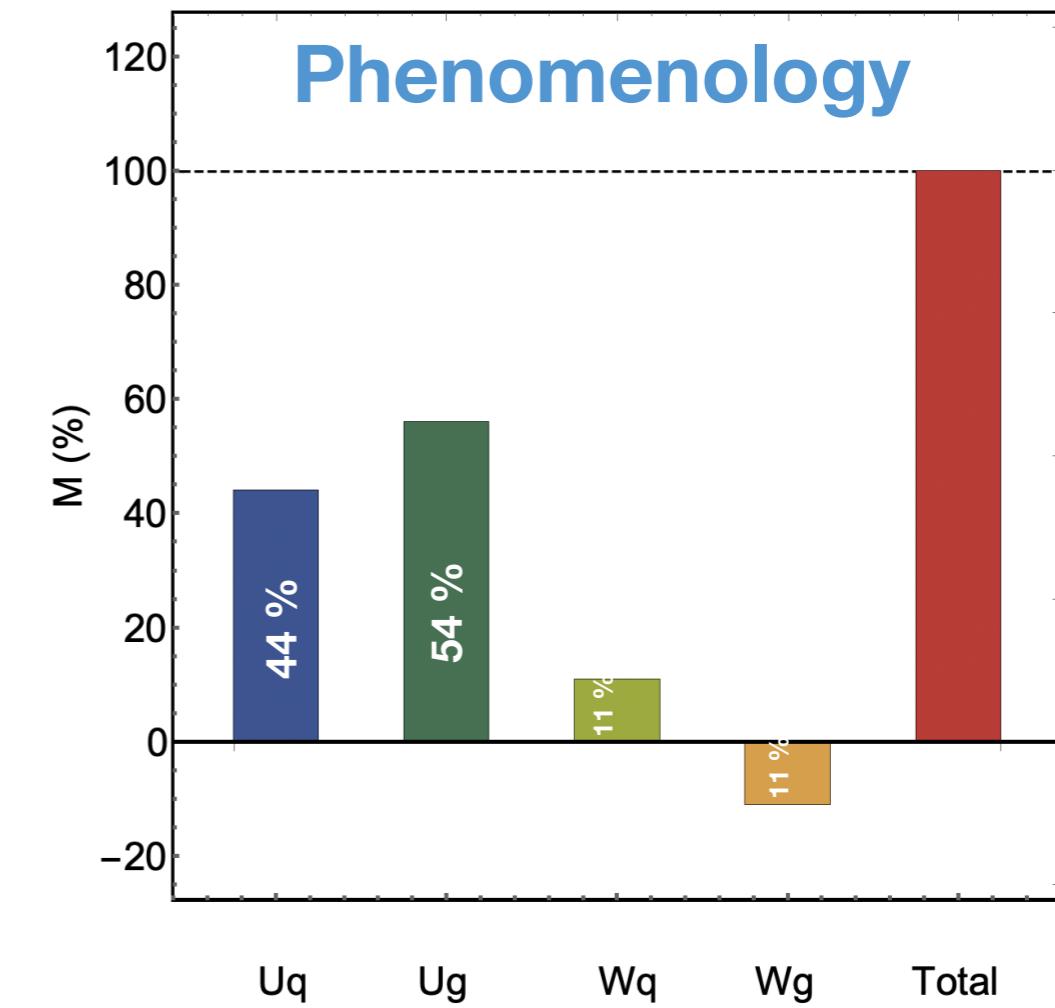
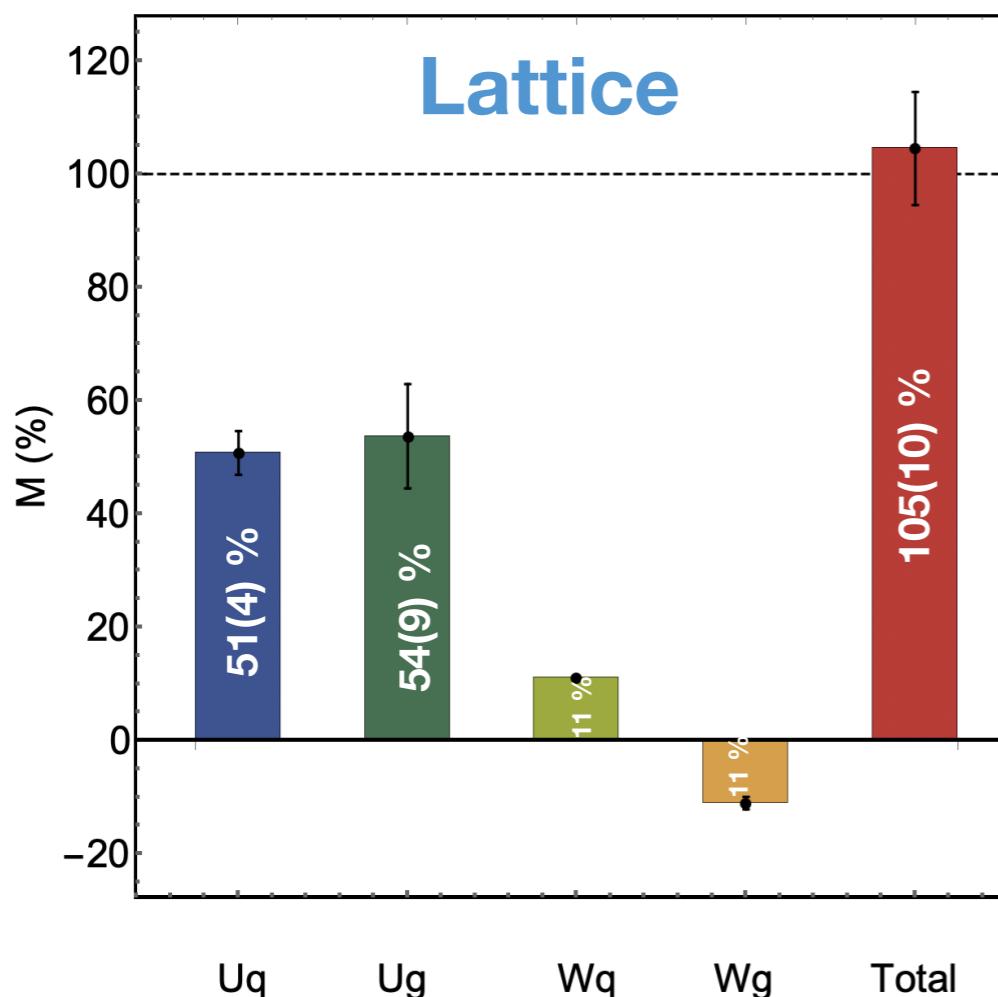
energy density and pressure kept separately

- ★ Ji's components ($T_{q,g}^{00} = \hat{T}_{q,g}^{00} + \bar{T}_{q,g}^{00}$) interpreted as internal energy
(effective coupled four-fluid picture)

$$\mathcal{W}_m = -M_m, \quad \mathcal{W}_q = \frac{1}{3} \left(M_q + \frac{12M_m}{4 + \gamma_m} \right), \quad \mathcal{W}_g = \frac{M_g}{3}, \quad \mathcal{W}_a = -M_a$$

Lorce's Pressure-volume work decomposition

two-fluid picture



- ★ Decomposition of lattice data gives the **same picture** as phenomenology

[C. Lorce', EPJ. C78 (2018) 2]
[L. Harland-Lang et al., EPJ. C 75 (2015)]
[M. Hoferichter et al., PRL 115 (2015)]

- ★ Equal contributions to the mass from the internal quark and gluon energies

- ★ U_g, W_g use sum rule:

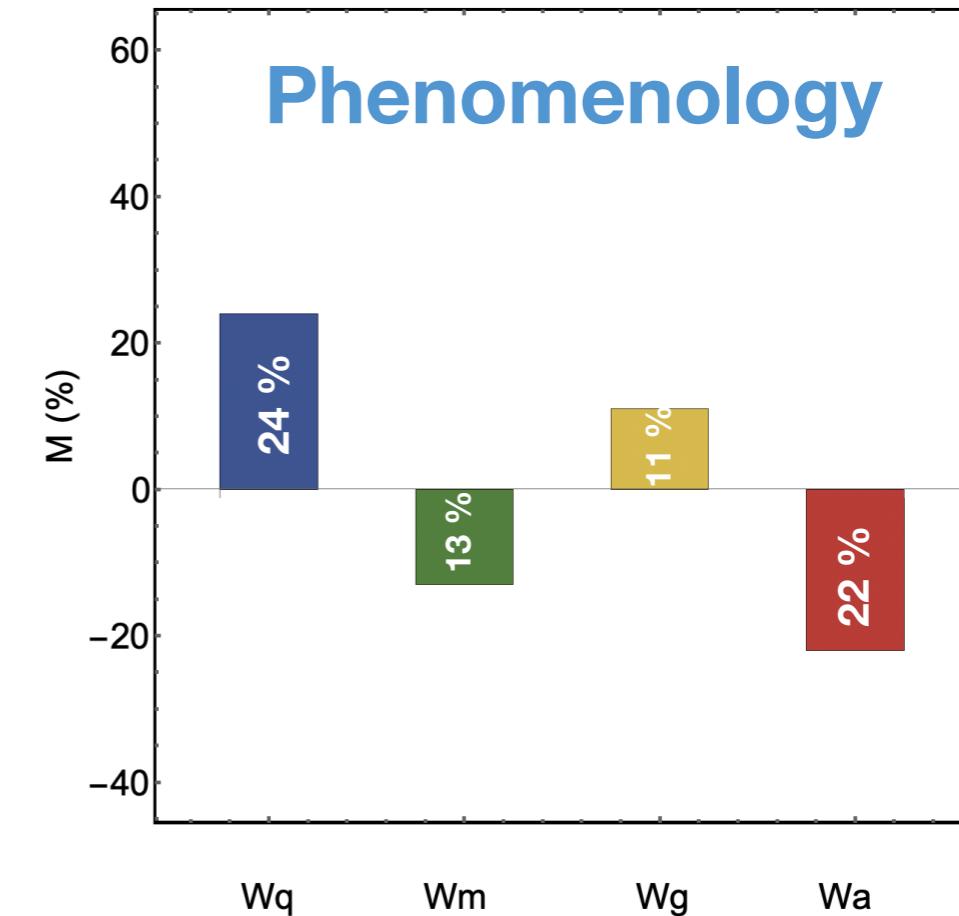
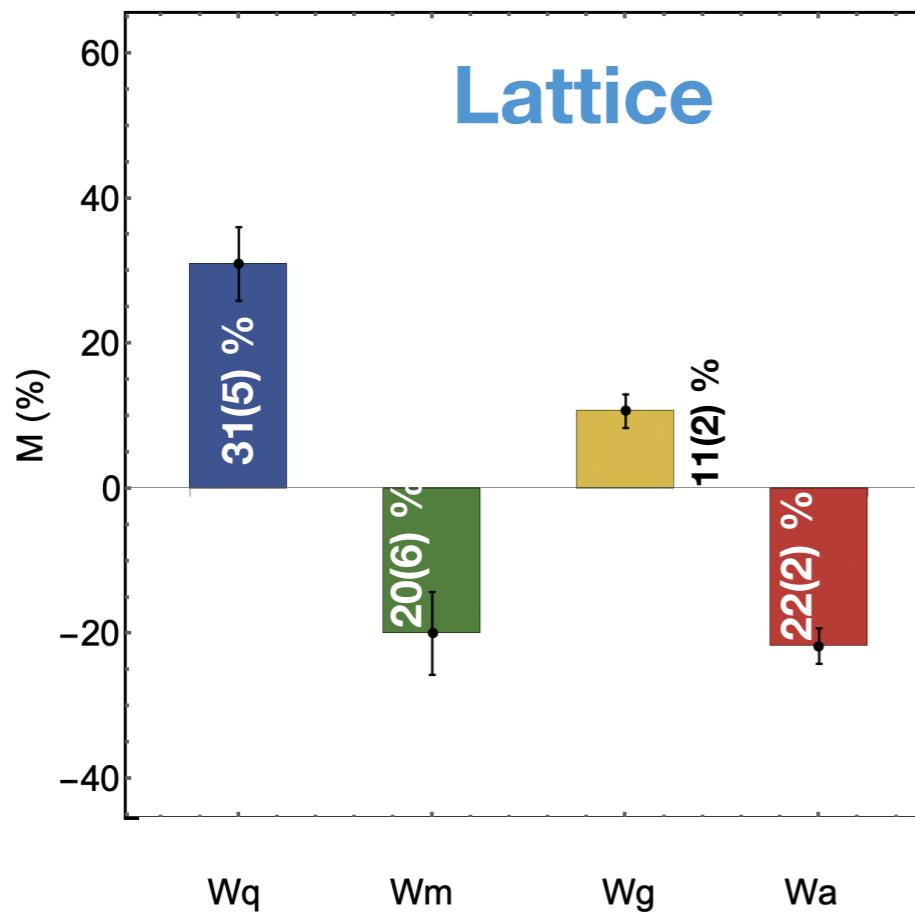
$$\bar{C}_g(0) + \bar{C}_q(0) = 0$$

		$O(\alpha_s^1)$	$O(\alpha_s^2)$	$O(\alpha_s^3)$
Scenario A	U_q	0.384 ± 0.035	0.383 ± 0.036	0.384 ± 0.036
	U_g	0.554 ± 0.035	0.556 ± 0.036	0.555 ± 0.036
Scenario B	U_q	0.420 ± 0.016	0.420 ± 0.017	0.421 ± 0.017
	U_g	0.518 ± 0.016	0.518 ± 0.017	0.517 ± 0.017

[A. Metz, B. Pasquini, S. Rodini, arXiv:2006.11171]

Ji's Pressure-volume work decomposition

four-fluid picture



- ★ Lattice and pheno data give similar picture
- ★ Total “quark” contribution (W_g , W_m) similar to total “gluon” contribution (W_a , W_g)
- ★ W_a , W_g receive input from sum rule
- ★ $\sum W_q = 0$ by construction

[C. Lorce', EPJ. C78 (2018) 2]
[L. Harland-Lang et al., EPJ. C 75 (2015)]
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Hatta-Rajan-Tanaka

Decomposition

$$\boxed{\frac{\langle T_\mu^\mu \rangle}{\langle N|N \rangle} = M}$$

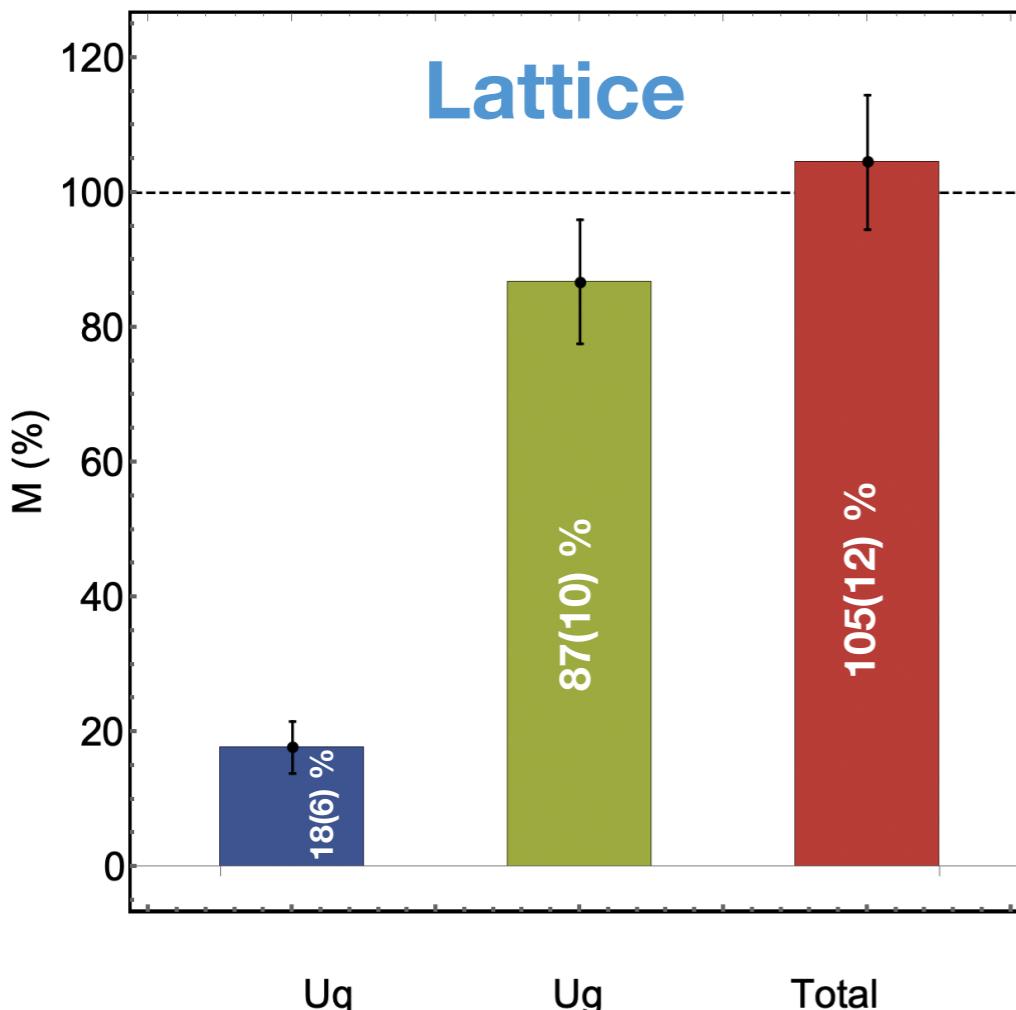
HRT Decomposition

★ [Y. Hatta, A. Rajan, K. Tanaka, JHEP 12, 008 (2018) arXiv:1810.05116; K. Tanaka, JHEP 01, 120 (2019), arXiv:1811.07879]

★ Separation of quark and gluon parts of trace part of EMT, $T_{\mu; q,g}^{\mu}$

$$T_{\mu; q}^{\mu} = (1 + c_1^{\text{MS}}) m \bar{\psi} \psi + c_2^{\text{MS}} F^{\alpha\beta} F_{\alpha\beta}$$

$$T_{\mu; q}^{\mu} = (\gamma_{\mu} - c_1^{\text{MS}}) m \bar{\psi} \psi + \left(\frac{\beta(g)}{2g} - c_2^{\text{MS}} \right) F^{\alpha\beta} F_{\alpha\beta}$$



$$U_{q,g} = M \left[A_{q,g}(0) + 4 \bar{C}_{q,g}(0) \right]$$

Sum rules

$$M = U_q + U_g$$

★ Gluon contributions is dominant
(in support of the argument that gluons are responsible for the mass)

Concluding Remarks

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“Can lattice calculate the mass distribution in the nucleon?”



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TMD Topical Collaboration

DOE Early Career Award (NP)
Grant No. DE-SC0020405



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“Can lattice calculate the mass distribution in the nucleon?”

- ★ Mass components are local 2-parton operators calculable in LQCD
- ★ Successful calculation of nucleon sigma terms, quark and gluon momentum fractions



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- ★ Successful calculation of nucleon sigma terms, quark and gluon momentum fractions

“Can one calculate the anomaly contribution on the lattice?”

- ★ Well-defined operator, but complicated renormalization pattern, and suppressed signal-to-noise ratio
- ★ Sum rules very useful to extract trace anomaly indirectly

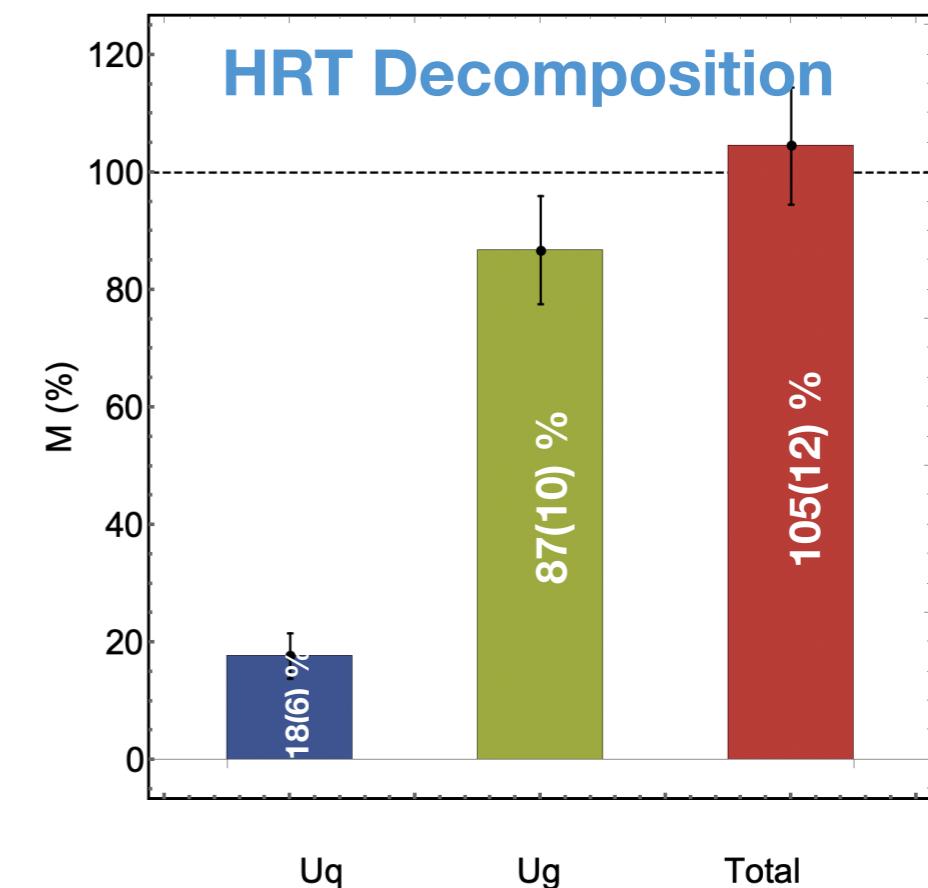
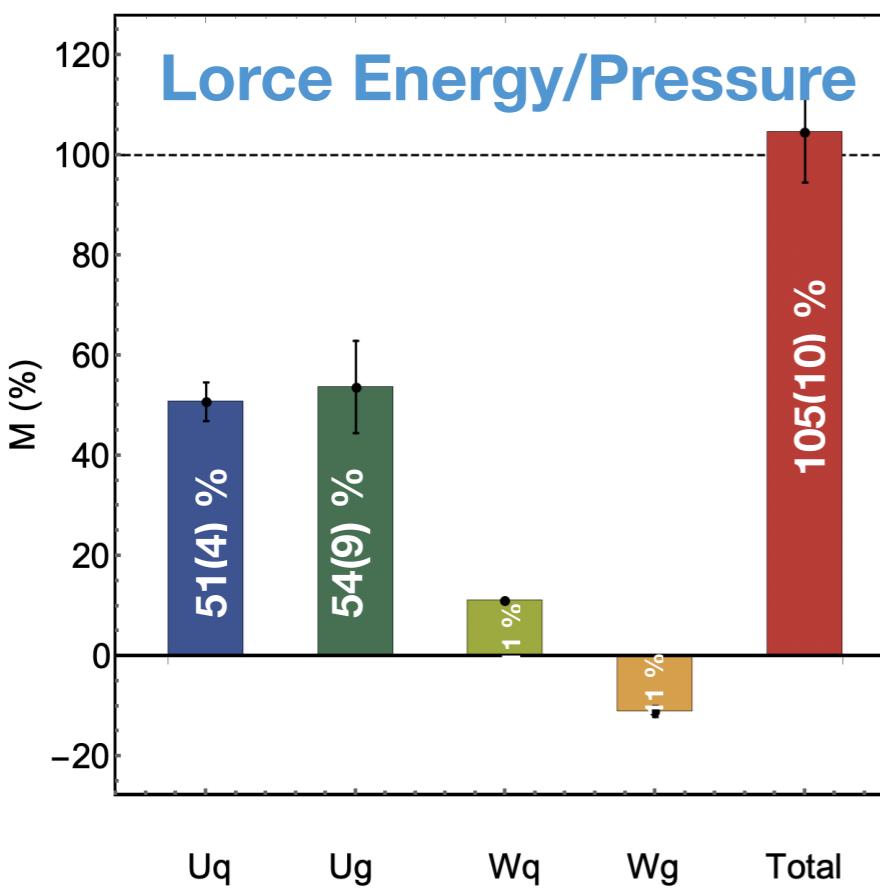
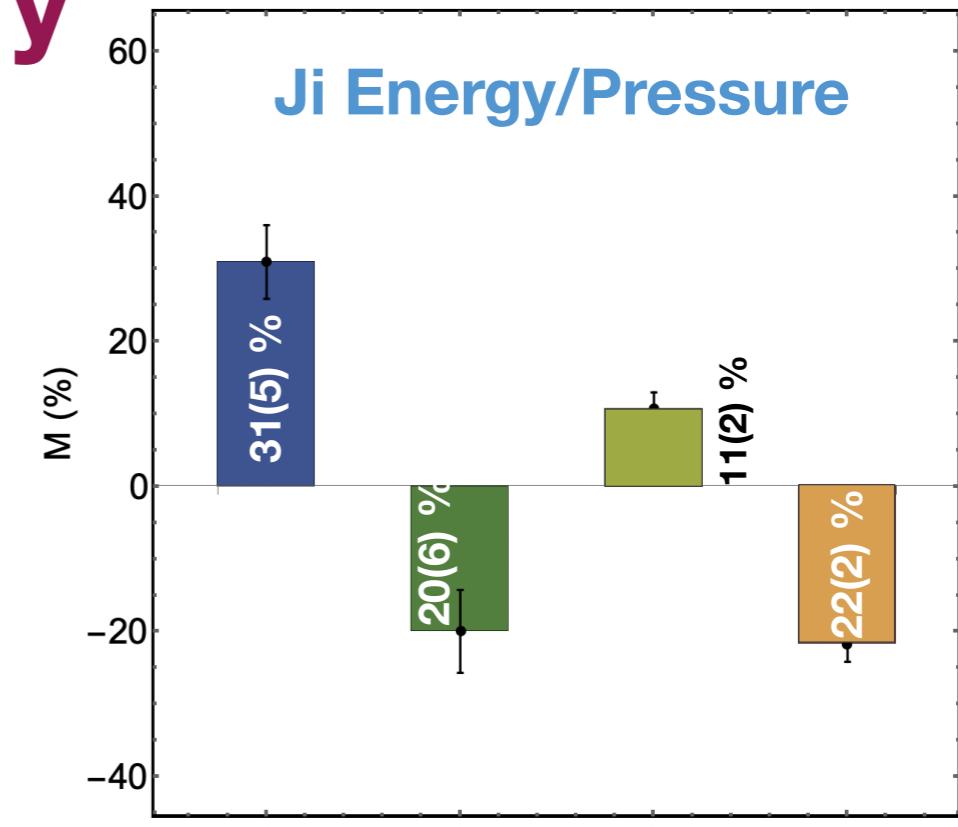
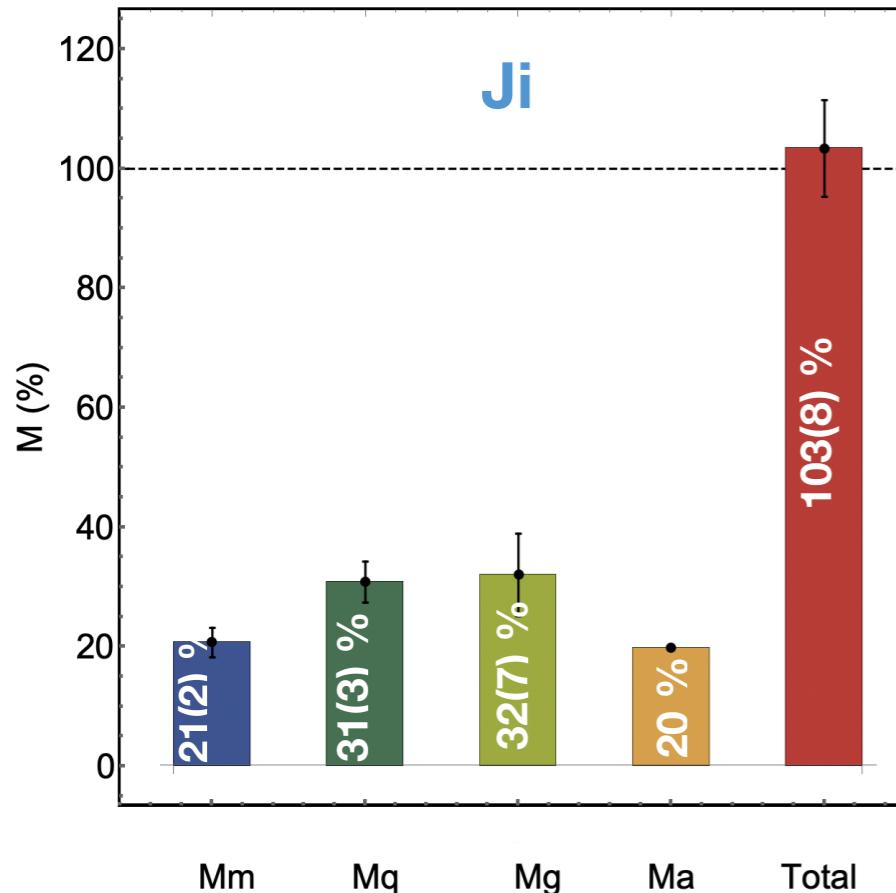


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Summary



Thank you